# High-energy X-ray Diffraction Study of Orbital Physics in $LnVO_3$ (Ln = La, Ce, and Y)

Y. Ren,<sup>1</sup> A.A. Nugroho,<sup>2</sup> T.T.M. Palstra,<sup>2</sup> J. Strempfer,<sup>3</sup> U. Rütt,<sup>3</sup> C.W. Kimball<sup>4</sup>

<sup>1</sup>Argonne National Laboratory, Argonne, IL, U.S.A.

<sup>2</sup>Solid State Chemistry Laboratory, Materials Science Centre, University of Groningen, Groningen, The Netherlands <sup>3</sup>Max-Planck-Institut für Festkörperforschung, Stuttgart, Germany

<sup>4</sup>Department of Physics, Northern Illinois University, DeKalb, IL, U.S.A.

## Introduction

In correlated electronic systems, the interplay between charge, spin, orbital degrees of freedom, and lattice distortion gives rise to a large variety of interesting physical properties (e.g., high-T<sub>c</sub> superconductivity, colossal magnetoresistance [CMR] effect, ferroelectrics). It has been recently realized that the orbital degree of freedom plays an important role in these phenomena. Because of its strong coupling with charge, spin, and lattice dynamics, it often manifests itself in various structural responses (e.g., charge stripes, diffuse charge/orbital-ordering scattering. induced-phase transition). Among these responses, the orbital physics in transition metal oxides is of particular interest [1]. Highenergy x-ray scattering is a powerful technique for probing the subtle structural changes that contain information that is essential for understanding the fundamental physics in these highly correlated systems.

The transition metal oxides with partially filled  $t_{2g}$ orbitals exhibit interesting phenomena, because of the relative weak Jahn-Teller coupling, high degeneracy, and greater orbital freedom when compared with the e<sub>g</sub> orbitals. Recently, we have found temperature-induced magnetization reversals and a novel phase transition in YVO<sub>3</sub> single crystals [2]. The transition involves competition between single-ion anisotropy and antisymmetric interaction and a change of two types of spin and orbital ordering in a noncubic symmetry [3], where the G-type spin-ordering (G-SO) is associated with the C-type orbital-ordering (C-OO), and the C-type spinordering (C-SO) is associated with the G-type orbitalordering (G-OO) (Fig. 1) [4]. However, in LaVO<sub>3</sub> and CeVO<sub>3</sub>, the situation is different because they form in a pseudo-cubic structure with almost equal V-V bonds above T<sub>N</sub>, in which quantum fluctuations among orbitals are expected to play an important role in the C-type magnetic ordering where both FM and AF coupling coexist. In addition to the relatively simple electronic configurations, it is a challenging task to understand the complex interactions between different degrees of freedom. It is thus of particular interest to perform detailed and systematic theoretical and experimental studies of the orbital physics in these systems.



FIG. 1. Two types of spin- and orbital-ordering and associated lattice distortion in LnVO<sub>3</sub>.

### Methods, Materials, and Results

The structural phase transition in the orthovanadates LaVO<sub>3</sub> and CeVO<sub>3</sub> has been studied with high-resolution, high-energy synchrotron x-ray diffraction at beamline station 11-ID-C at the APS. LaVO<sub>3</sub> undergoes a secondorder phase transition at  $T_N = 143$ K and a first-order transition at  $T_t = 141$ K (Fig. 2). In CeVO<sub>3</sub>, however, there are second-order phase transitions occurring at  $T_0 = 154$ K and first-order phase transitions occurring at  $T_N = 134$ K. These phase transitions are confirmed by specific heat measurements. The phase transition at  $T_t$  in LaVO<sub>3</sub> or  $T_0$ in  $CeVO_3$  is due to a *G*-type orbital ordering, which lowers the structure symmetry from orthorhombic Pbnm to monoclinic  $P2_1/b11$ . The structure change at  $T_N$  in CeVO<sub>3</sub> is ascribed to an orbital-ordering-enhanced magnetostrictive distortion. The change at  $T_N$  in LaVO<sub>3</sub>, however, is most probably due to an ordered occupation of the vanadium  $3d t_{2g}$  orbitals associated with antiferromagnetic ordering. We propose that the



FIG. 2. Lattice parameters versus temperature of LaVO<sub>3</sub>.

first-order phase transition at  $T_t$  in LaVO<sub>3</sub> should be associated with a sudden change of both the spin and orbital configurations, similar to the phase transition at  $T_s = 77$ K in YVO<sub>3</sub> [2-4], causing a reversal of the net magnetization. However, the ordered state above  $T_t$  in LaVO<sub>3</sub> is identical to that below  $T_s$  in YVO<sub>3</sub>. It has been found that with a decreasing lanthanide ionic radius from La, the Neel temperature  $T_N$  decreases while the orbital ordering onset temperature increases in these orthovanadates (Fig. 3). However, when the  $Ln^{3+}$  size further decreases, the system shows much more complex behavior, which is the subject of our current investigations.

#### Acknowledgments

Use of the APS was supported by the U.S. Department of Energy, Office of Science, Office of Basic Energy



FIG. 3. Magnetic-  $(T_N)$  and orbital-ordering temperature  $(T_o)$  versus Ln ionic radius.

Sciences, under Contract No. W-31-109-ENG-38. See Y. Ren, A.A. Nugroho, A.A. Menovsky, J. Strempfer, U. Rütt, F. Iga, T. Takabatake, and C.W. Kimball, "Orbital-ordering-induced phase transition in LaVO<sub>3</sub> and CeVO<sub>3</sub>," Phys. Rev. B **67**, 014107 (January 2003) for more information.

#### References

[1] Y. Tokura and N. Nagaosa, Science **288**, 462-468 (2000).

[2] Y. Ren, T.T.M. Palstra, D.I. Khomskii, E. Pellegrin, A.A. Nugroho, A.A. Menovsky, and G.A. Sawatzky, Nature **396**, 441-444 (1998).

[3] Y. Ren, T.T.M. Palstra, D.I. Khomskii, A.A. Nugroho, A.A. Menovsky, and G.A. Sawatzky, Phys. Rev. B **62**, 6577-6586 (2000).

[4] G.R. Blake, T.T.M. Palstra, Y. Ren, A.A. Nugroho, A.A. Menovsky, Phys. Rev. Lett. **87**, 245501 (2002).