# Pressure Effect on Structure and Conductivity of Sc- and Al-Tungstates

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### Introduction

Because of high electrostatic interaction with the surrounding framework structure of anions such as  $O^{2-}$ , highly charged trivalent cations are often considered to be poor candidates for ionic migration. Recently, it was reported that trivalent cationic conduction is possible in a  $Sc_2(WO_4)_3$ -type structure [1, 2]. This structure consists of a framework that is linked by corner-sharing  $ScO_6$  octahedra and  $WO_4$  tetrahedra and that offers a large tunnel size for ionic migration. In this study, we investigated the effects of high pressure (P) on the structure and conductivity of Sc-and Al-tungstates as a means of exploring a new route to high-conductivity behavior.

## **Methods and Materials**

 $Sc_2(WO_4)_3$  and  $Al_2(WO_4)_3$  were synthesized by standard solid-state techniques from the constituent oxides. Measurements of ac conductivity in a frequency range of 10-1 MHz at 400°C and at pressures up to 4.3 GPa were carried out in a 1000-ton cubic press by using a Solartron SI 1260 impedance/gain-phase analyzer. High-P, *in situ* x-ray diffraction (XRD) experiments were performed in a 250-ton DIA-type high-P apparatus at beamline 13-BM-D at the APS. A NaCl pressure standard was used for *in situ* measurement of P. Energy dispersive XRD spectra were collected at each P setting with a Ge detector at a fixed 2qangle.

## **Results and Discussion**

P-dependence of  $Sc_2(WO_4)_3$  conductivity at 400°C is plotted in Fig. 1 and compared with the 1-atm value [2]. Relative to the 1-atm value, the high-P data show a decreasing trend with increasing P up to ~2.9 GPa, where the conductivity dependence on P reverses. Conductivity increases with increasing P up to the highest P of 4.3 GPa. This latter behavior is unexpected for ionic conduction. With decreasing P, conductivity decreases montonically, with a final trajectory in the low-P region toward the 1-atm value.

In search of a structural explanation of the positive P-dependence of conductivity above 2.9 GPa, high-P XRD patterns for  $Sc_2(WO_4)_3$  at 400°C were taken at several pressures up to 4.4 GPa (Fig. 2). At the detector 2*q* angle of 3.82°, the main diffraction peaks from  $Sc_2(WO_4)_3$  are localized between 40 and 50 keV. Upon pressurization, the



FIG. 1. Dependency of  $Sc_2(WO_4)_3$  conductivity on pressure at 400°C.



FIG. 2. Pressure dependency of  $Sc_2(WO_4)_3$  conductivity at 400°C compared with 1-atm value.

sharp peaks present at ambient conditions quickly broaden, intensities drastically decrease, and the peaks shift to higher energy up to 2.2 GPa. The diffraction peak intensities decrease again with a further increase in P up to 4.4 GPa, when all diffraction peaks from the sample are absent.

The decrease in peak intensity and broadening in peak shape observed for  $Sc_2(WO_4)_3$  with increasing pressure is interpreted as gradual P-induced amorphization. The  $Sc_2(WO_4)_3$  belongs to a relatively small class of substances with a negative thermal expansion coefficient [3, 4]. The structure of Sc<sub>2</sub>(WO<sub>4</sub>)<sub>3</sub> consists of ScO<sub>6</sub> octahedra and WO<sub>4</sub> tetrahedra linked through corner-sharing to form a 3-D framework [5, 6]. The negative value of expansion is related to transverse vibrations of bridging oxygens in the Sc-O-W linkages [3, 4]. Anharmonicity in these vibrations leads to cooperative tilting of the quasi-rigid framework polyhedra, causing an effective decrease in the average Sc-O-W bridging bond angle and a concomitant densification upon heating. The complex coupled twisting in three dimensions of the relatively rigid framework polyhedra upon heating provides a possible route to P-induced disorder [7]. The overall 3-D framework structure can perhaps be best understood by breaking it down into 2-D slabs of cornersharing polyhedra parallel to the a-c plane, which stack together to form the overall structure. The c and a axes are more compressible than the b axis on the basis of the flexible hinged Sc-O octahedra and W-O tetrahedra linkages. The Sc-O-W angles in the polyhedra within the *a*-*c* plane are affected the most by P. This leads to the polyhedra exhibiting a higher degree of anisotropy in the a-c plane relative to either the b-c or the a-b plane. This degree of anisotropy increases with compression, until the rigid polyhedra begin to oppose each other. The forces that tend to keep polyhedra regular are mainly the oxygen-oxygen repulsive interactions within the polyhedron. To accommodate further volume reduction as P is increased, the polyhedra must rotate or tilt and eventually rupture the flexible Sc-O-W linkages. This is macroscopically identified as amorphization. The conductivity results at P above 2.9 GPa suggest that amorphization enhances conduction. We interpret this to be the result of the increase in structural entropy and its concomitant effects on accessibility of neighboring sites to mobile  $Sc^{3+}$ .

The association of pressure-induced amorphization behavior and negative thermal expansion behavior, as well as the unexpected conductivity behavior in  $Sc_2(WO_4)_3$ , led us to investigate the pressure behavior of  $Al_2(WO_4)_3$  which is iso-structural with  $Sc_2(WO_4)_3$  but has a positive thermal expansion behavior. The results of high-P XRD showed that the structure remains crystalline up to the maximum experimental pressure. The pressure dependences of the lattice parameters *a*, *c*, calculated from XRD patterns at 400°C, are shown in Fig. 3. All parameters exhibit a normal decrease with an increase in pressure below 2.5 GPa. However, an obvious change occurs at pressures between 2.8 and 3.2 GPa for all lattice parameters. The two short crystallographic axes, *a* and *c*, of the orthorhombic unit cell approach one another in this critical pressure region, and



FIG. 3. Pressure dependencies of lattice parameters a, c calculated from XRD patterns at 400°C.

they remain approximately equal upon any further increase of pressure. This signals a pressure-induced structural transition from orthorhombic to tetragonal. In the 3-D  $Sc_2(WO_4)_3$ -type framework structure, there are densely populated corner-sharing layers of WO<sub>4</sub> tetrahedra and  $ScO_6$ octahedra. These layers extend parallel to the *a*-*c* plane and stack together along the *b* direction only through O(3) to form the overall structure [3]. Such a stacking results in the *b* axis being more compressible than the *a* and *c* axes at pressure.

Conductivity data (Fig. 4) also show clearly that  $Al_2(WO_4)_3$  exhibits very different behavior at high P than does  $Sc_2(WO_4)_3$ . Conductivity increases very fast with increasing pressure below 2.0 GPa, then slowly increases and reaches a maximum value of ~ $1.6 \times 10^{-2}$  (ohm · cm<sup>-1</sup>), which is about five orders of magnitude greater than the 1-atm value, at pressures close to 4.0 GPa. Upon decreasing pressure, conductivity remains constant until a pressure of 0.5 GPa, when the sample cooled down from 400°C to room temperature. Since the dramatic change in conductivity occurs below 2.0 GPa, it cannot be attributed to the structural change at 2.5 GPa. The measurement of dc polarization reveals that electrical conduction in  $Al_2(WO_4)_3$  has been changed from ionic dominant to electronic dominant at high P.

Mixed valence states of  $W^{6+}$  and  $W^{5+}$  have been observed in the bronzes  $[A_x(P_2O_4)_2(WO_3)_{2m}]$ ,  $Na_xWO_3$ , and other crystalline and glass tungsten oxides [9-11]. The presence of mixed-valent W causes a delocalization of the electrons over the polyhedral framework. These materials exhibit metallic, electronically conducting properties. The concentration of the lower valence state,  $W^{5+}$ , and therefore the conductivity, is strongly influenced by octahedral ion substitution or the environmental atmosphere [10, 11]. The high-P closed environment, with the sample encapsulated in BN, is a reducing atmosphere that is oxygen deficient. Therefore, the



FIG. 4. Conductivity data.

change in ionic to electronic conduction can then be attributed to the tungsten valence change in  $Al_2(WO_4)_3$  at high P, which results in electron transfer through  $W^{5+} \rightarrow W^{6+}$ .

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