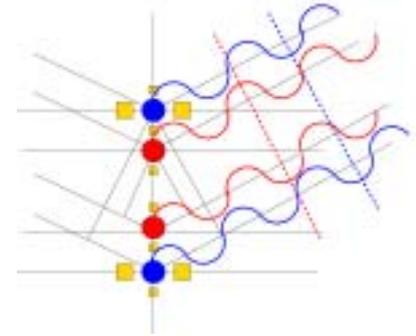


Beyond element-specific magnetism: Magnetic spectroscopy in the diffraction channel

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Advanced Photon Source

Nanomagnetism using x-ray techniques
Lake Geneva, WI, August 2004



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* LDRD funding

Argonne National Laboratory



A U.S. Department of Energy
Office of Science Laboratory
Operated by The University of Chicago



Outline

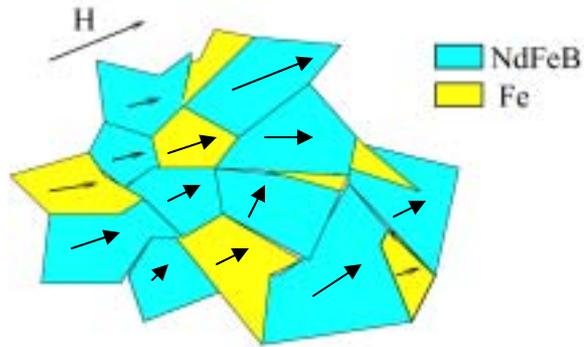
- The Need: to go beyond element specific magnetism.
- The Tool: Resonant diffraction of CP x-rays.
- Example: Site-specific magnetism in $\text{Nd}_2\text{Fe}_{14}\text{B}$.
- Magnetic spectroscopy in the diffraction channel.
- Outlook and Summary.

The Need: to go beyond element specific magnetism

Element-specific XMCD lacks phase, site selectivity.

Phase-Specific

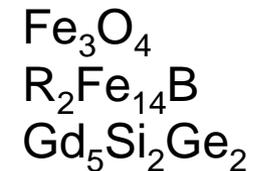
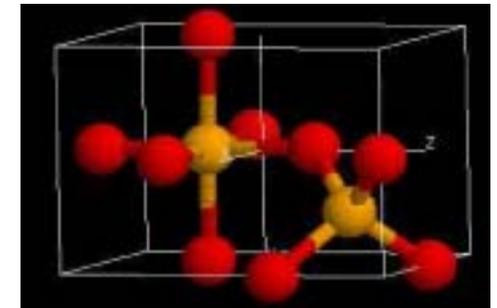
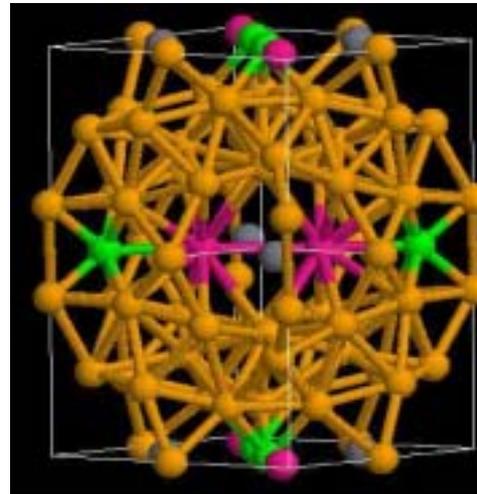
Same element in coexisting phases



Nanocomposites: NdFeB/Fe
Impurity phases: $\text{Sm}_2\text{Co}_7/\text{SmCo}_5$
Intrinsic (CMR, stripes)

Site-specific

Same element in inequivalent crystal sites

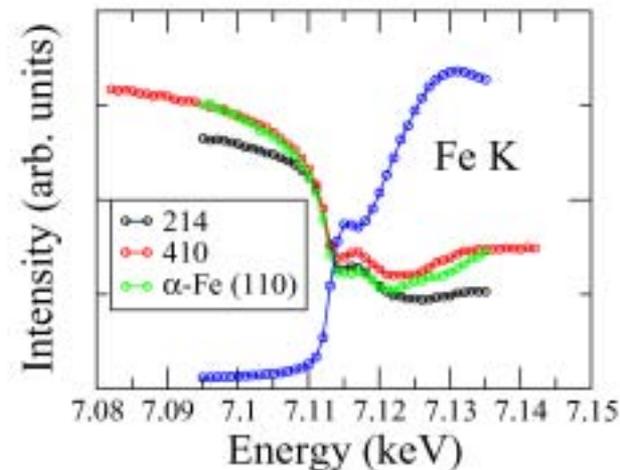
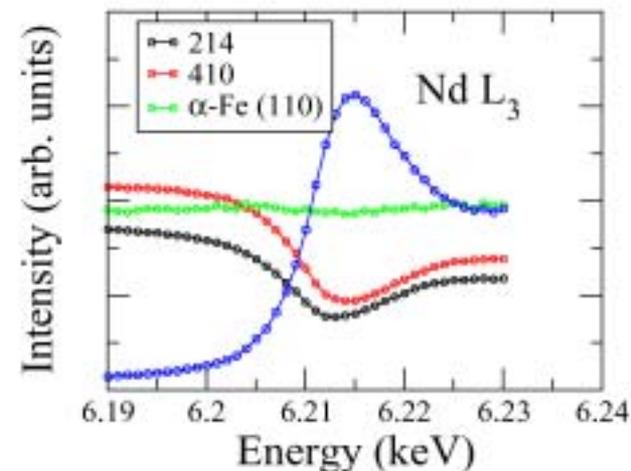
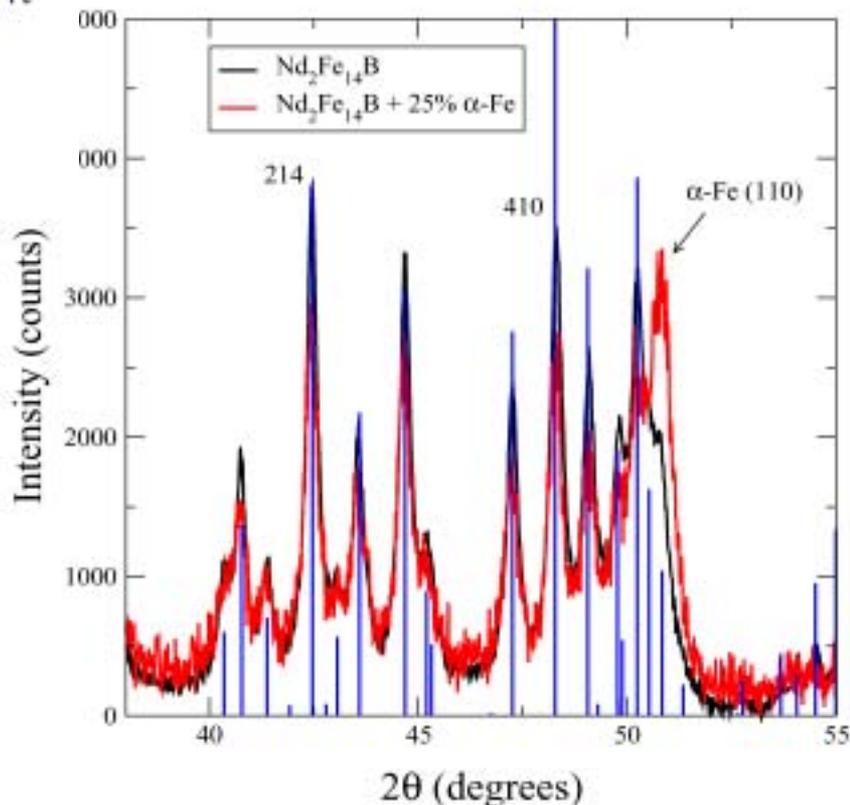
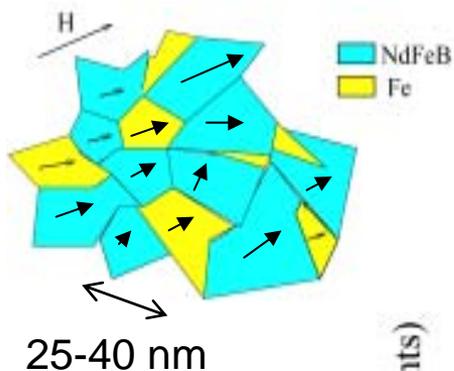


The Tool: resonant diffraction of CP x-rays

Diffraction = phase and site selectivity

Atomic resonance = element specificity

CP x-rays = coupling to sample magnetization



The Tool: resonant diffraction of CP x-rays

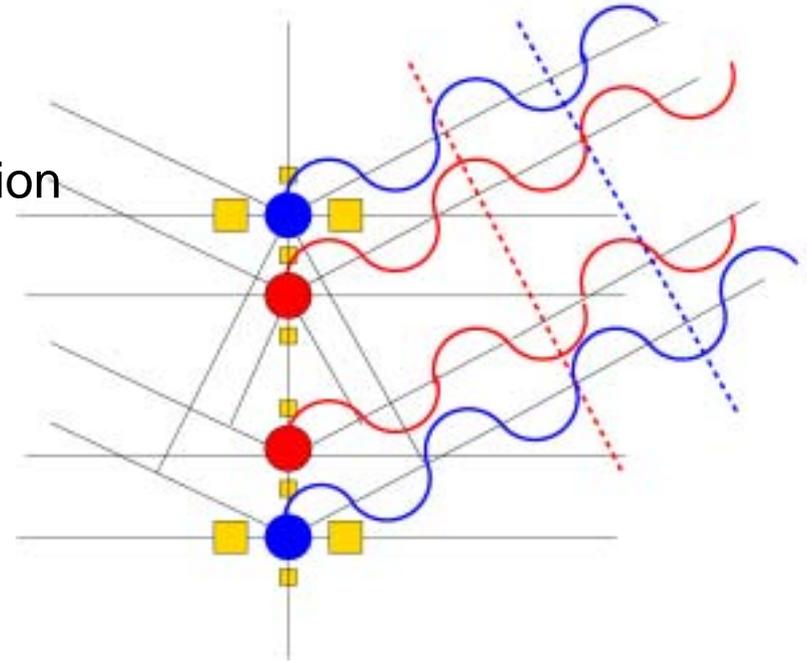
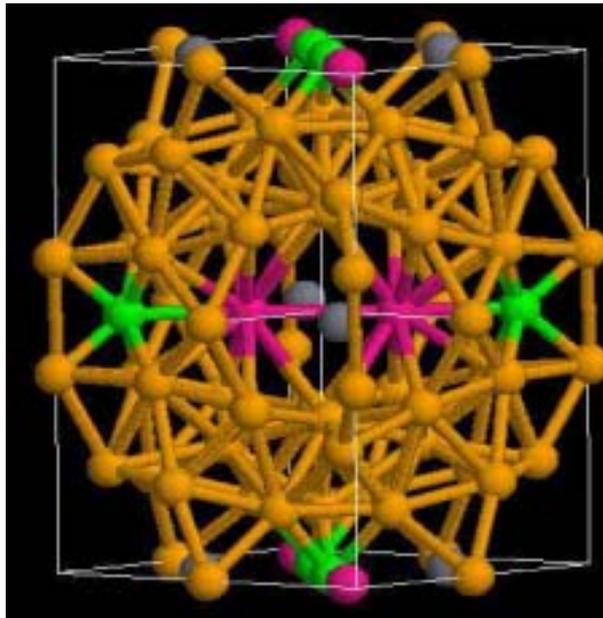
Diffraction = phase and site selectivity

Atomic resonance = element specificity

CP x-rays = coupling to sample magnetization



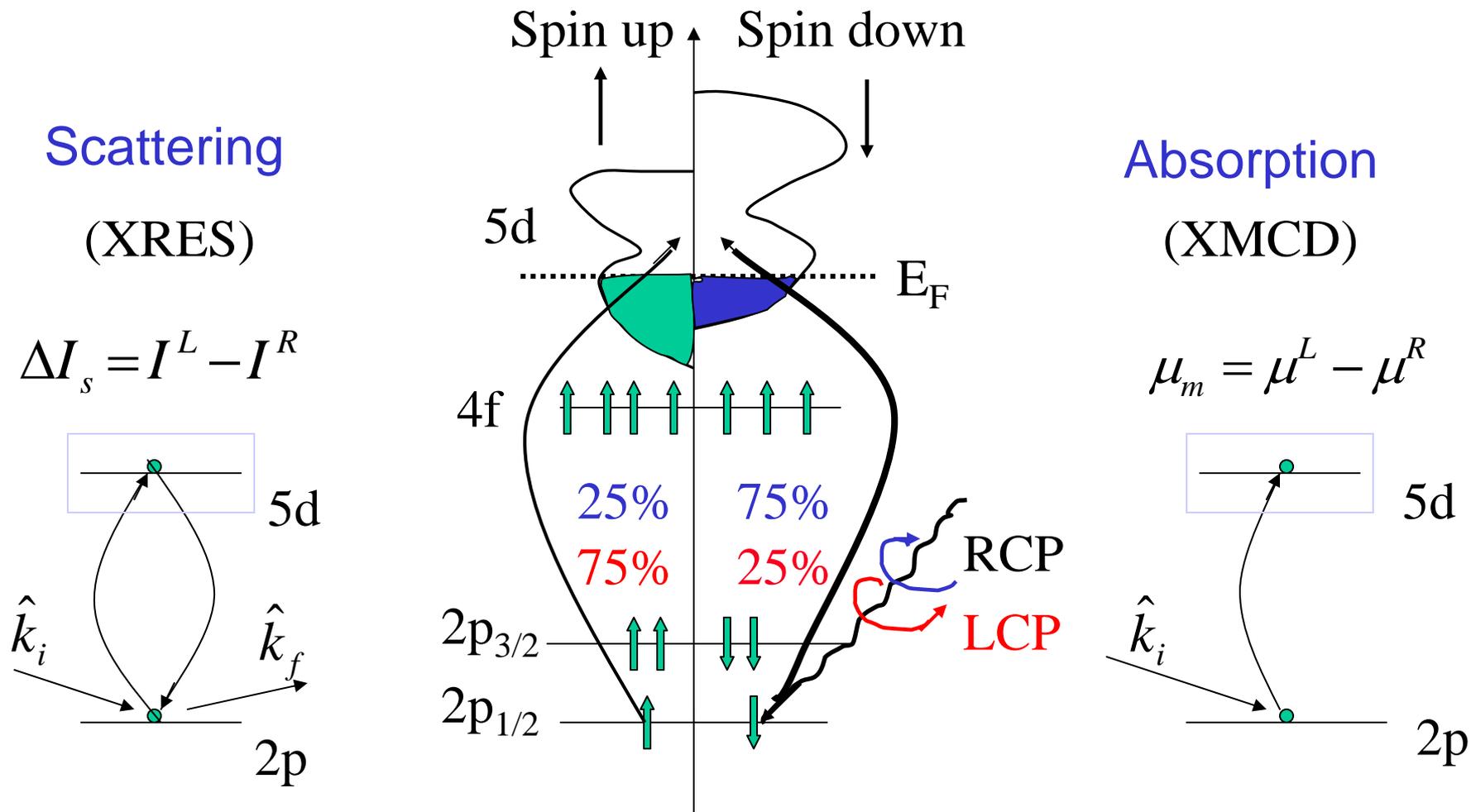
- Nd (4g)
- Nd (4f)
- B
- Fe



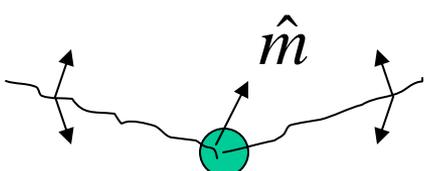
$$(4g) (x, \bar{x}, 0); (\bar{x}, x, 0); \left(\frac{1}{2} + x, \frac{1}{2} + x, \frac{1}{2}\right); \left(\frac{1}{2} - x, \frac{1}{2} - x, \frac{1}{2}\right)$$

$$(4f) (x, x, 0); (\bar{x}, \bar{x}, 0); \left(\frac{1}{2} + x, \frac{1}{2} - x, \frac{1}{2}\right); \left(\frac{1}{2} - x, \frac{1}{2} + x, \frac{1}{2}\right)$$

Resonant absorption and scattering of CP x-rays



Scat. Amp. (E1) for magnetic ion near resonance (Hannon *et al.*, 1988)

$$f = (f_0 + f_e)(\epsilon'^* \cdot \epsilon) + if_m (\epsilon'^* \times \epsilon) \cdot \hat{m} + \cancel{f^{mag}}$$


Charge
1

Resonant
charge
0.2-0.5

Resonant
Magnetic
0.01-0.5

Non-resonant magnetic
0.001

$$I \propto f^2 = \text{charge} + \text{magnetic} + \underline{\text{charge} \times \text{magnetic}}$$

Needs to get rid of strong charge scattering (ferro-ferrimagnets)

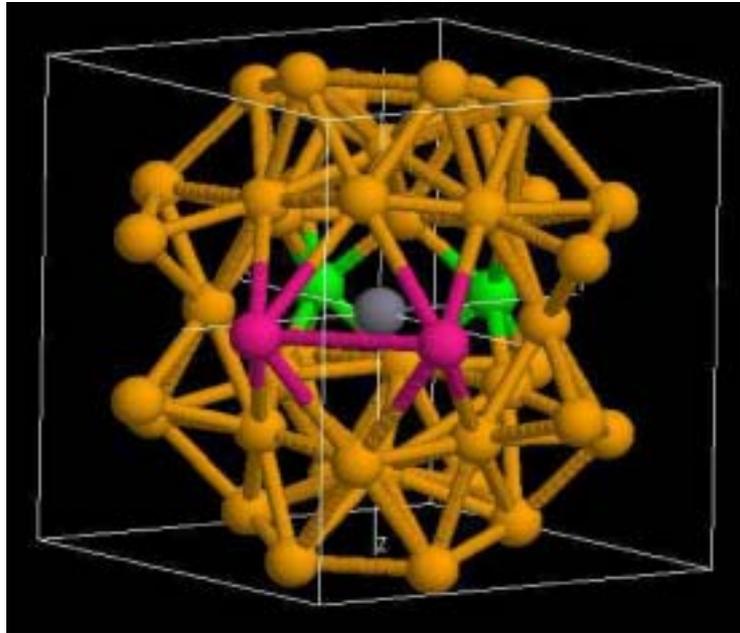
$$\Delta I_s = I^L - I^R \propto \{ i [(\epsilon'^* \cdot \epsilon)^* \underbrace{(\epsilon'^* \times \epsilon) \cdot \hat{m}}_{\hat{m} \parallel \text{scat. plane}}] \underbrace{(f_0 + f_e)^*}_{\text{Strong E-dependence}} \cdot \underbrace{f_m}_{\text{Strong E-dependence}} \} + c.c.$$

$\hat{m} \parallel \text{scat. plane}$

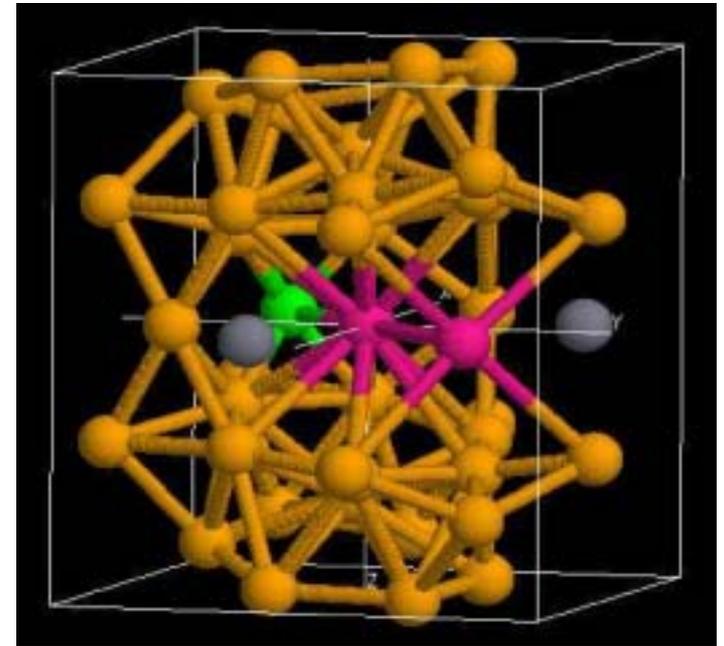
Strong E-dependence

$$\Sigma I_s = I^L + I^R = A(f_0 + f_e)^2 + \cancel{Bf_m^2}$$

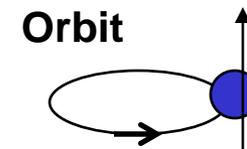
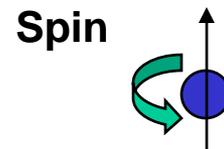
Example: Site-specific magnetism in $\text{Nd}_2\text{Fe}_{14}\text{B}$



- Nd 4g
- Nd 4f
- B
- Fe



- Fe: $\sim 31 \mu_{\text{B}}/\text{f.u.}$
- Nd: $\sim 6 \mu_{\text{B}}/\text{f.u.}$

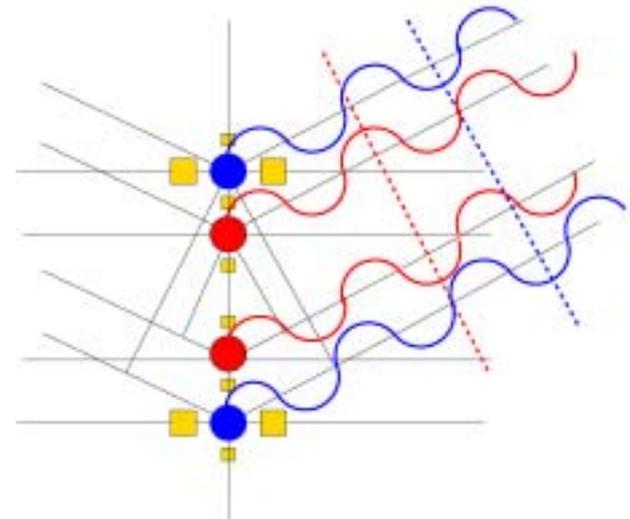


- Easy axis: $[001]$ at RT, tilts towards $[110]$ at low T
- Magnetic “hardness” due to large orbital moment of Nd ions.

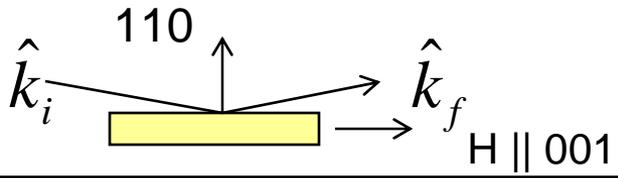
Site-specific Nd structure factors in Nd₂Fe₁₄B

$$F_{hkl} = f_{Nd}(Q, E) \sum_{i=1}^4 e^{i\vec{Q} \cdot \vec{r}_i}$$

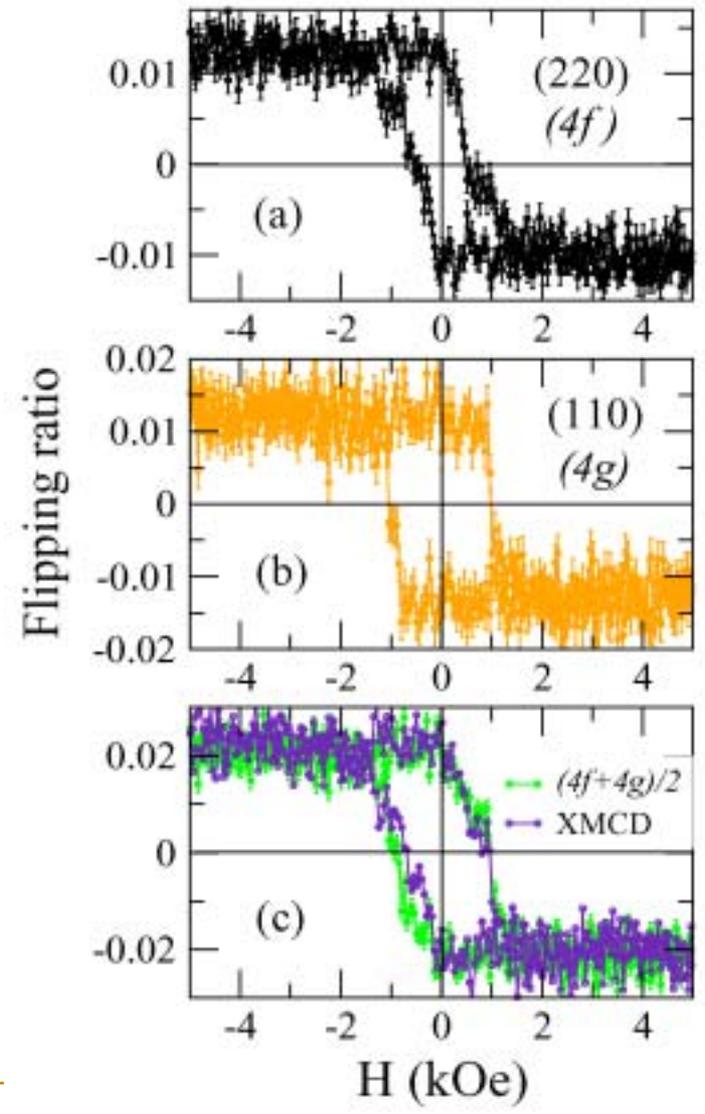
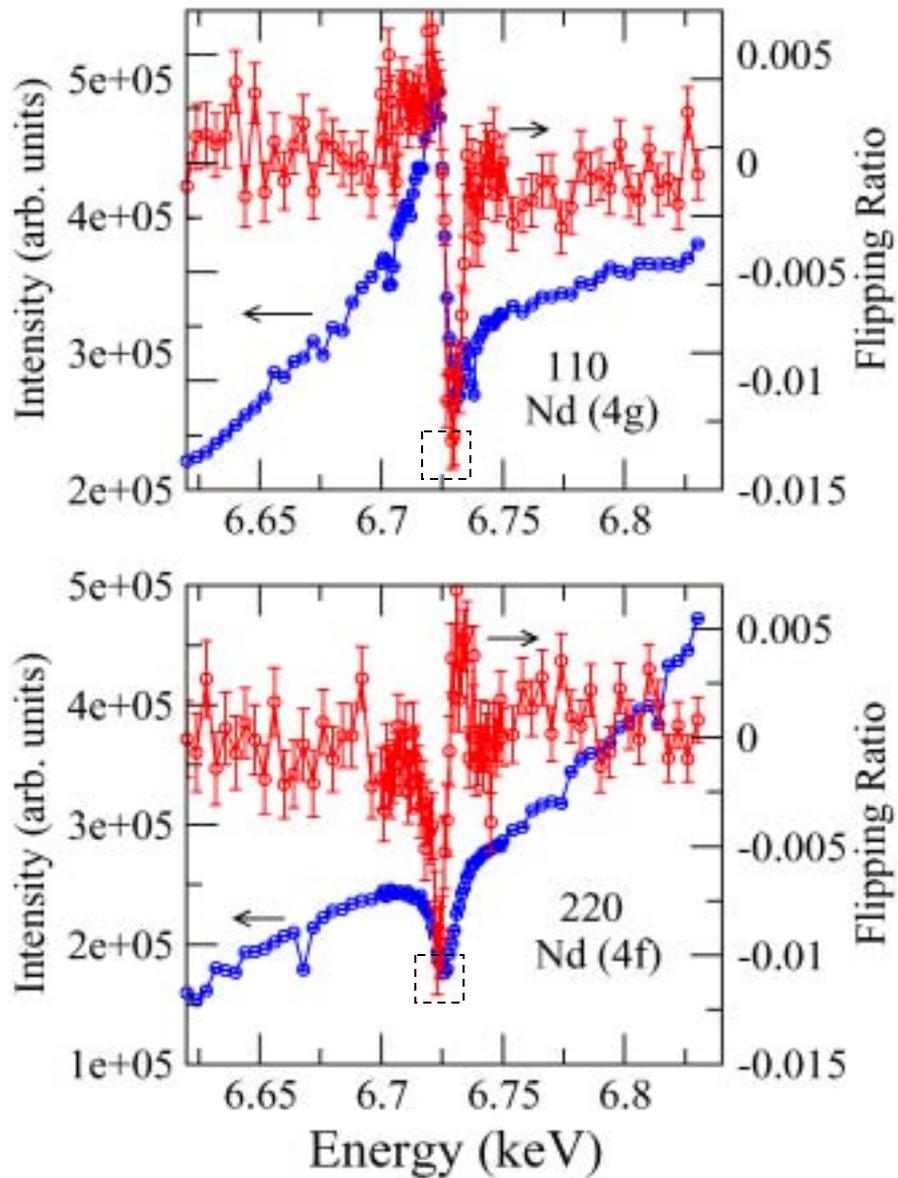
$$\vec{Q} = 2\pi \left(\frac{h}{a}, \frac{k}{b}, \frac{l}{c} \right), \vec{r}_i = (x_i a + y_i b + z_i c)$$



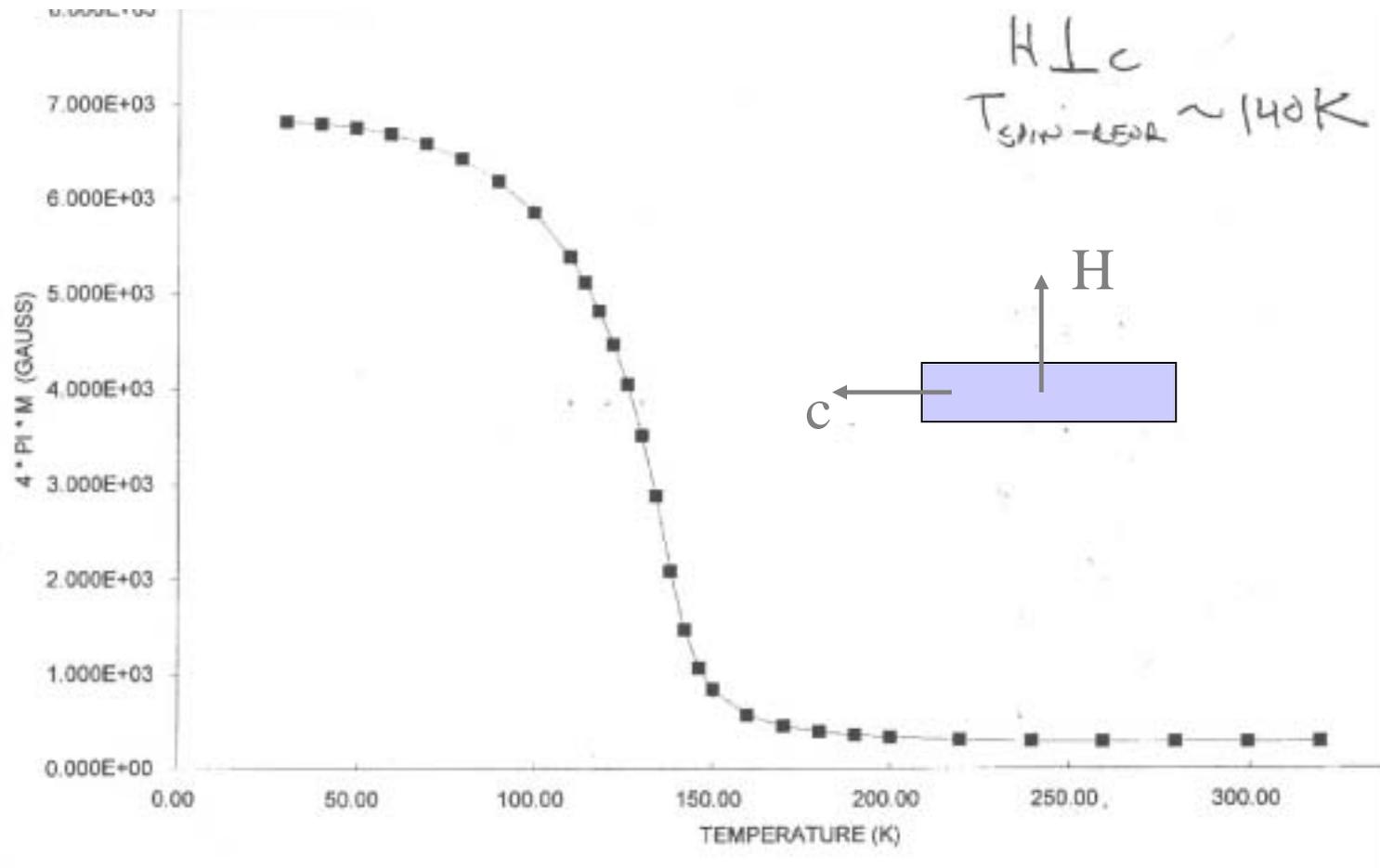
Wyckoff Site	(110)	(220)	(440)
4f	2.2+0.5i	140+35i	86+29i
4g	69+15i	5.2+1.3i	92+31i



$$I = |f_e(\boldsymbol{\varepsilon}'^* \cdot \boldsymbol{\varepsilon}) + if_m(\boldsymbol{\varepsilon}'^* \times \boldsymbol{\varepsilon}) \cdot \hat{m}|^2$$



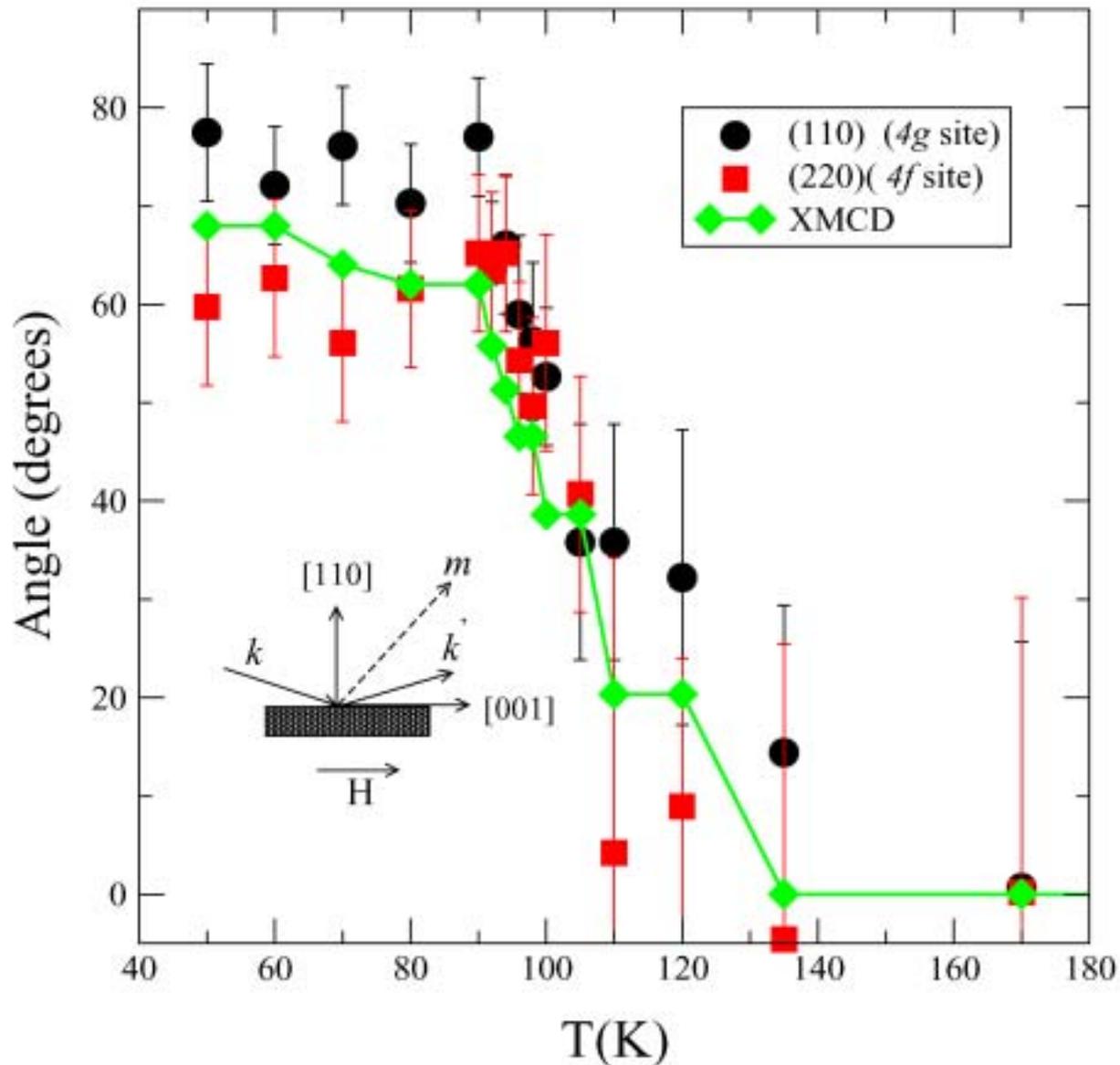
Spin reorientation transition in $\text{Nd}_2\text{Fe}_{14}\text{B}$



Paul Canfield, Ames Lab



Site-specific spin reorientation transition



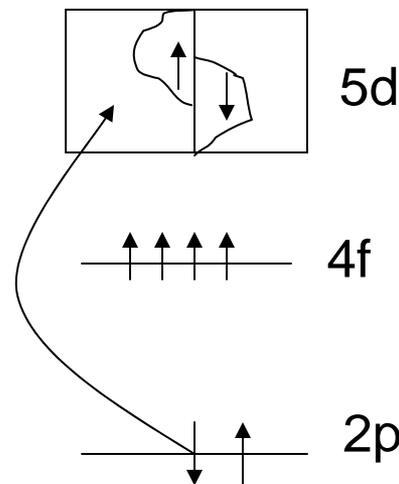
Magnetic spectroscopy in the diffraction channel

- Spin-dependent DOS (E_F).
- Size of moment (orbital, spin).
- Orientation of moment, magnetization reversal.

Easy interpretation in absorption channel (XMCD)
[optical theorem $\mu_m \propto f''_m(q=0)$].

$$\mu_m \propto N(\uparrow) - N(\downarrow)$$

$$\mu_m \propto M \cos \theta$$

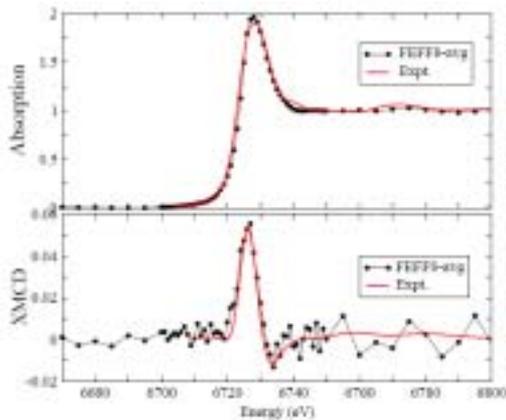
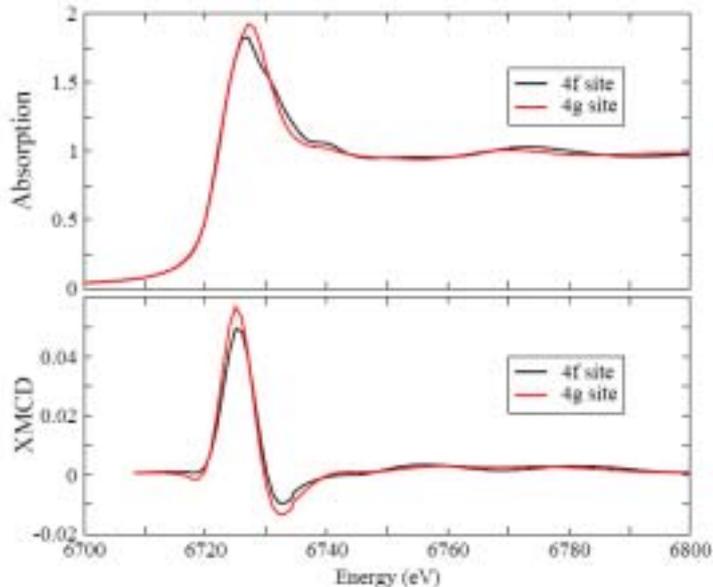


Resonant diffraction can yield same information, but harder to retrieve.
(Intensity depends on real and imaginary parts of f_e , f_m ; phase problem).

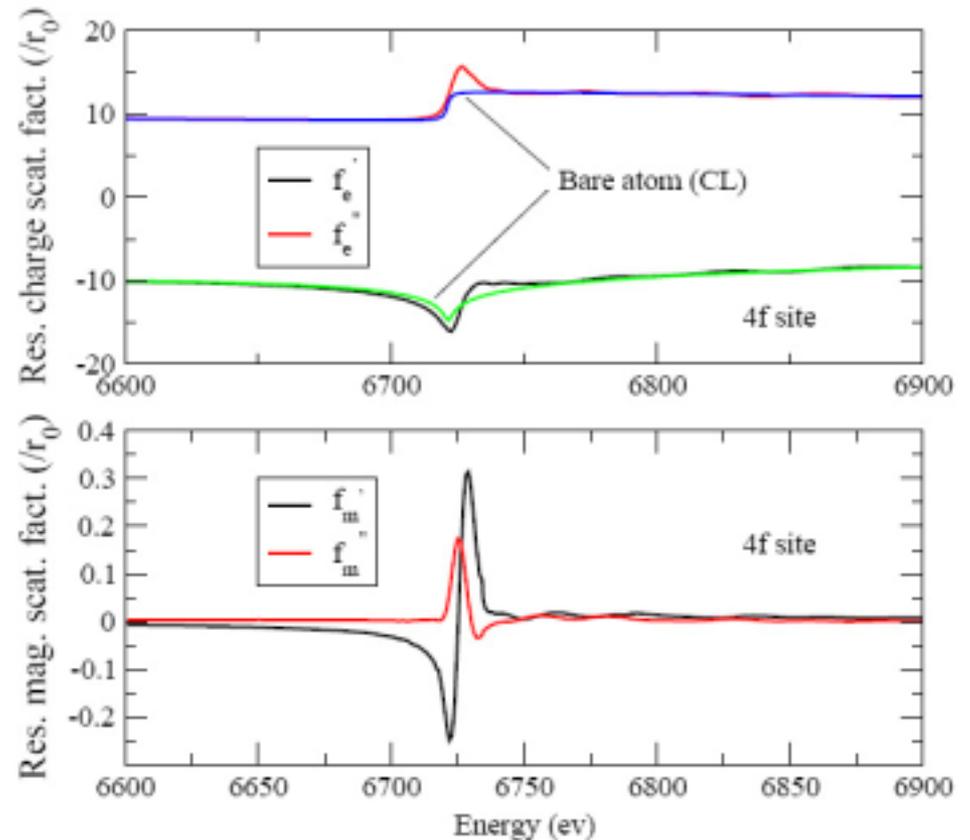


Magnetic spectroscopy in the diffraction channel

Ab-initio (FEFF8) calculations of μ_e , μ_m for each site separately.



$$\mu(E) \rightarrow f''(E) \rightarrow f'(E)$$



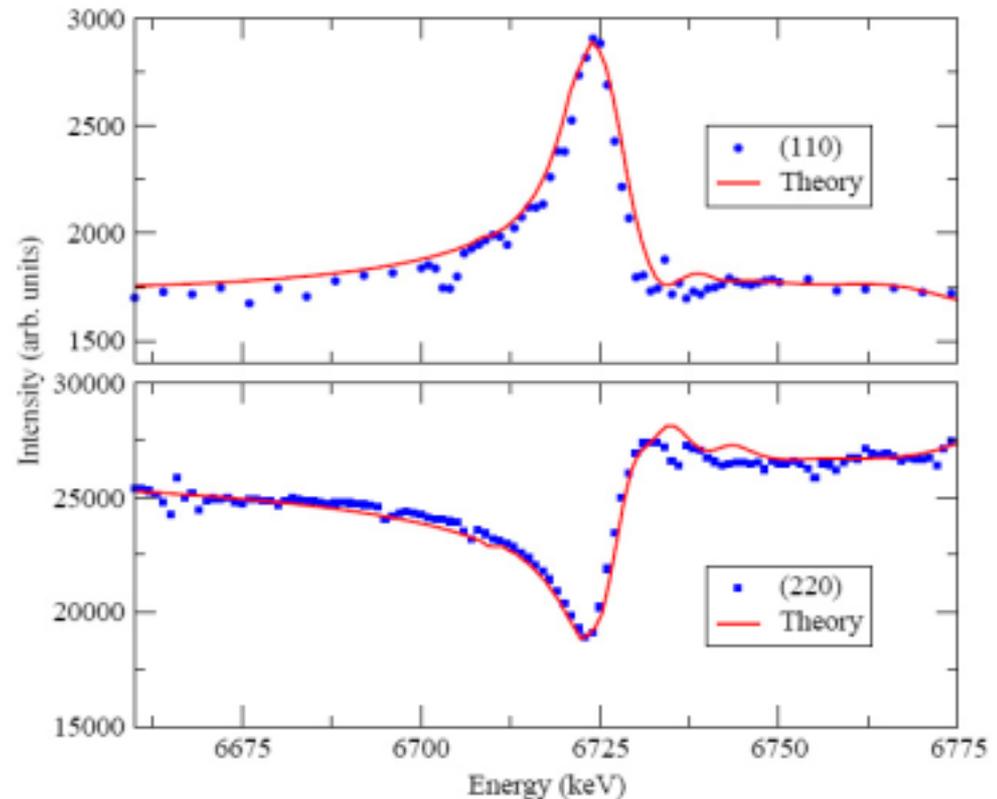
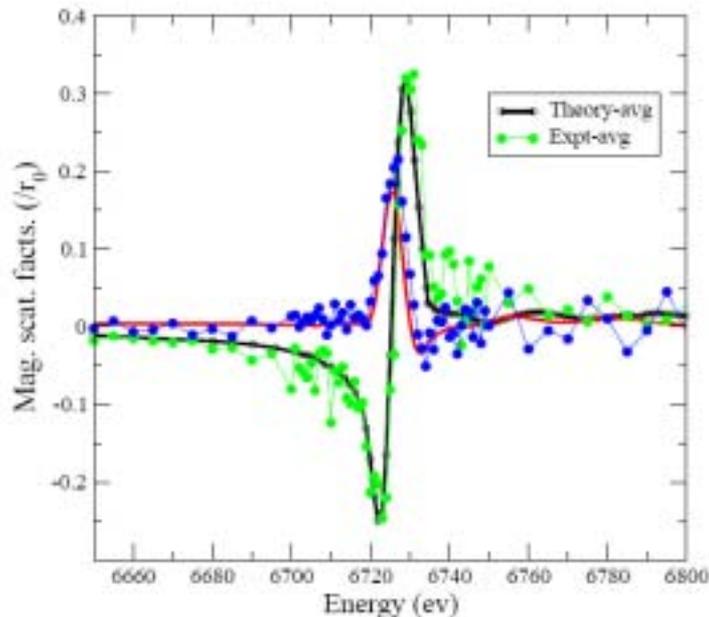
Magnetic spectroscopy in the diffraction channel

$$F_q = \sum_i (A f_{e,i} + i B f_{m,i}) e^{i\vec{q} \cdot \vec{r}_i}$$

$$\Delta I_s = 2i \sum_{i,j} (-A B^* f_{m,i}^* f_{e,j} + A^* B f_{e,i} f_{m,j}^*) e^{i\vec{q} \cdot (\vec{r}_i - \vec{r}_j)}$$

$$\Sigma I_s = 2 \sum_{i,j} A^2 f_{e,i}^* f_{e,j} e^{i\vec{q} \cdot (\vec{r}_i - \vec{r}_j)}$$

Structure factor calcs with ab-initio res. scat. factors as initial guess; refine.



Outlook and summary

- Diffraction (**Q**) + resonance (**E**) + CP x-rays (**M**) allows for phase and site specific studies of magnetism in artificially or naturally inhomogeneous systems.
- Combination with ab-initio calculations should yield phase/site specific spectroscopic information, such as SDOS(E_F) and magnetic moments.

